- 5. Ya. S. Podstrigach and Yu. M. Kolyano, "Equations of generalized thermoelasticity for bodies with thin inclusions," Dokl. Akad. Nauk SSSR, <u>224</u>, No. 4, 794-797 (1975).
- Ya. S. Podstrigach and Yu. M. Kolyano, "Heat exchange taken into account in local heating of thin-walled structural elements," Dokl. Akad. Nauk SSSR, <u>225</u>, No. 4, 778-781 (1975).

CALCULATION OF TOMOGRAPHIC PROJECTIONS

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The article suggests a method of calculating tomographic projections.

The problem of interaction between x rays and the substance of the investigated object, arising in the field of computerized tomography, reduces to the calculation of tomographic projections [1]. The present article submits a method of calculating parallel and bundle tomographic projections for one class of images of the section of the object; the terminology and some of the designations are taken over from [1].

Let w and \hat{w} be the applicates of points of the plane of the object's section in the initial system x, y and in the system of coordinates \hat{x} , \hat{y} rotated through the angle θ , respectively, $\hat{w} = w e^{-i\theta}$; let $\mu(x, y)$ and $\mu_{\theta}(x, \hat{y})$ be the distribution of the absorption coefficient by the material of the object in the initial and in the rotated system of coordinates, respectively, $\mu_{\theta}(\hat{x}, \hat{y}) = \mu(x, y)$; the function $\mu(x, y)$ is called the image of the section of the object. Then for x rays passing along the straight line $\hat{x} = \text{const}$, the logarithm of the ratio of its intensity at the entrance into the object to the intensity at the exit from the object, called the parallel tomographic projection $\rho_{\theta}(\hat{x})$ of the section, is determined by the formula

$$p_{\theta}(\hat{x}) = \int_{-\infty}^{+\infty} \mu_{\theta}(\hat{x}, \hat{y}) d\hat{y}.$$
(1)

Assume that from the source lying at the point $\left[i\left(\beta - \frac{1}{2}\right)\right]$ there emerges a beam in the direction parallel to the vector exp $\left[i\left(\frac{\pi}{2}+\beta+\gamma\right)\right]$; the logarithm of the ratio of its intensities at the entrance into and at the exit from the object is called the bundle projection $h_{\beta}(\gamma)$ of the section; it is correlated with the parallel projection by the relation [1]

$$h_{\beta}(\gamma) = p_{\theta(\beta,\gamma)}(x(\beta,\gamma)), \tag{2}$$

where the dependences $\hat{x}(\beta, \gamma), \theta(\beta, \gamma)$ have the form

$$\hat{x} = -\rho \sin \gamma, \ \theta = \beta + \gamma. \tag{3}$$

We introduce the notation: l, n are integers, n = 1, 2, ..., N; $l = 1, 2, ..., L_n$; g(n, l) is the region bounded by an ellipse with the center at the point $R(n, l) \exp[i\varphi(n, l)]$, the semi-axis a(n, l), b(n, l), the first of which is inclined to the radius vector of the center of the ellipse at the angle $\Phi(n, l)$; if for some n_0, l_0 we have R $(n_0, l_0) = 0$, then we put $\varphi(n_0, l_0) = 0$; g(0) is the region bounded by an ellipse with the center at the origin of coordinates, semiaxes a(0), b(0), the first of which is inclined to the x axis at the angle $\Phi(0)$.

Let us examine the class of images $\mu(\mathbf{x}, \mathbf{y})$ for which the following condition is fulfilled; the section of the object is the domain g(0); $g(n, l) \subset g(0)$ for all n, l; the sets G(n), determined by the relation

$$G(n) = \bigcup_{l=1}^{L_n} g(n, l),$$
 (4)

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do not intersect pairwise;

$$\mu(x, y) = \begin{cases} k(0), & w \in g(0) \setminus \bigcup_{n=1}^{N} G(n), \\ k(n), & w \in G(n), \end{cases}$$
(5)

where k(0) and k(n) are specified constants.

For any image $\mu(\textbf{x}, \, \textbf{y})$ from the examined class, the parallel projections are expressed by the formula

$$p_{\theta}(\hat{x}) = k(0) p_{\theta}(0, \hat{x}) - \sum_{n=1}^{N} [k(0) - k(n)] p_{\theta}(n, \hat{x}), \qquad (6)$$

the bundle projections are found from this with the aid of relations (2), (3). In formula (6) the value of $p_{\theta}(0, \hat{x})$ is determined by the equality

$$p_{\theta}(0, \hat{x}) = 2 \sqrt{A(0) a(0) b(0) - A^2(0) \hat{x}^2}, \qquad (7)$$

where

$$A(0) = \frac{\sin 2\alpha (0)}{1 - \cos 2\alpha (0) \cos 2 [\Phi(0) - \theta]}, \qquad (8)$$

$$\alpha(0) = \operatorname{arctg} \frac{a(0)}{b(0)} , \qquad (9)$$

and the values of $p_{\theta}(n, x)$ are calculated in the following manner.

Specifying the values n, \hat{x} , θ , and putting $l = 1, 2, ..., L_n$; j = -1, 1, we calculate $\hat{y}(n, l, j; \hat{x}, \theta)$ for all combinations of l, j by the formula

$$\hat{y}(n, l, j; \hat{x}, \theta) = R(n, l) \sin [\varphi(n, l) - \theta] + B(n, l; \theta) \times \\ \times C(n, l; \hat{x}, \theta) + (-1)^{\frac{3+j}{2}} \sqrt{A(n, l; \theta)a(n, l)b(n, l) - A^2(n, l; \theta)C^2(n, l; \hat{x}, \theta)},$$
(10)

where

$$A(n, l; \theta) = \frac{\sin 2\alpha(n, l)}{1 - \cos 2\alpha(n, l) \cos 2[\varphi(n, l) + \Phi(n, l) - \theta]}, \qquad (11)$$

$$B(n, l; \theta) = \frac{\cos 2\alpha (n, l) \sin [\varphi(n, l) + \Phi(n, l) - \theta]}{1 - \cos 2\alpha (n, l) \cos 2 [\varphi(n, l) + \Phi(n, l) - \theta]},$$
(12)

$$C(n, l; \hat{x}, \theta) = R(n, l) \cos [\varphi(n, l) - \theta] - \hat{x},$$
(13)

$$\alpha(n, l) = \operatorname{arctg} \frac{a(n, l)}{b(n, l)} .$$
(14)

We arrange the obtained set of values $\hat{y}(n, l, j; \hat{x}, \theta)$ (see (10)) in the form of a nondecreasing sequence of s, and denote:

$$\hat{y}(n, l_s, j_s; x, \theta) = y(s), s = 1, 2, \dots, S.$$
 (15)

The subscripts s for which the equality

$$\sum_{k=1}^{s} j_k = 0$$
 (16)

is fulfilled are arranged in the form of an increasing sequence of q: q = 1, 2, ..., Q. Then the sought value of $p_{\theta}(n, \hat{x})$ is determined by the expression

$$p_{\theta}(n, \hat{x}) = \sum_{q=1}^{Q} y(s_q) - \sum_{q=1}^{Q-1} y(s_q+1) - y(1), \qquad (17)$$

in the same way with the aid of formula (6) we find the parallel projection, and with a view to formulas (2), (3) we find the bundle projection for any image $\mu(x, y)$ from among the class under examination.

NOTATION

x, y, Cartesian coordinates; i, imaginary unit; r, φ , polar coordinates; $w = x + iy = r \exp(i\varphi)$; $\hat{w} = \hat{x} + i\hat{y} = w \exp(-i\theta)$; $\mu(x, y)$, absorption coefficient of radiation as a function of the coordinates (image of the section of the object); β , γ , bundle coordinates; $p_{\theta}(\hat{x})$, parallel tomographic projection; $h_{\beta}(\gamma)$, bundle tomographic projection.

LITERATURE CITED

1. H. J. Scudder, "Introduction to computer aided tomography," Proc. IEEE, <u>66</u>, No. 6, 628-637 (1978).

SURVEYS

INTERACTIONS OF ATOMS AND CALCULATION OF TRANSPORT COEFFICIENTS IN

METAL VAPORS AND THEIR MIXTURES WITH GASES

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In kinetic theory [1, 2], the transport coefficients of a gas are expressed in terms of collision integrals. In particular, in the first approximation of the theory of Chapman and Enskog, the viscosity and thermal conductivity of a dilute, single-component gas, and the coefficient of diffusion of a dilute binary mixture are given by the formula

$$\eta = \frac{5}{16} \frac{\sqrt{\pi m k T}}{\pi \sigma^2 \Omega^{(2,2)*}} , \qquad (1)$$

 $\lambda = \frac{5}{2} \eta c_v \tag{2}$

(for a monatomic gas),

$$D_{12} = \frac{3}{16nm_{12}} \frac{\sqrt{2\pi m_{12}kT}}{\pi \sigma_{12}^2 \Omega_{12}^{(1,1)*}} .$$
(3)

The reduced collision integrals

$$\Omega^{(l,s)*} = \Omega^{(l,s)} \left\{ \left(\frac{kT}{\pi m_{12}} \right)^{1/2} \frac{(s+1)!}{2} \left[1 - \frac{1 + (-1)^l}{2(l+1)} \right] \pi \sigma_{12}^2 \right\}^{-1}$$

are computed with the help of relations taking into account the interaction of molecules in collisions based on the conservation laws of mass, momentum, and kinetic energy:

$$\chi(g, b) = \pi - 2b \int_{R_m}^{\infty} \frac{dR/R^2}{\left[1 - b^2/R^2 - \varphi(R)/(m_{12}g^2/2)\right]^{1/2}},$$

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